Stochastic gradient descent and Random Matrix Theory

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with Chanju Park and Biagio Lucini

PRE 111 (2025) 1, 015303 [2407.16427 [cond-mat.dis-nn]]

and Ouraman Hajizadeh

2411.13512 [cond-mat.dis-nn]



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Random Matrix Theory (RMT)

developed by Wigner and Dyson to describe nuclear spectra (1959-1962)

o universal features: level spacing, Coulomb repulsion, Wigner surmise, fluctuations

non-universal behaviour: spectral density

example:

• successfully applied in QCD to describe Dirac operator

JJM Verbaarschot and T Wettig, *Random matrix theory and chiral symmetry in QCD* Ann. Rev. Nucl. Part. Sci. 50 (2000) 343 [arXiv:hep-ph/0003017 [hep-ph]].



RMT and machine learning

- Input layer Multiple hidden layers Output layer
- different context: machine learning and weight matrix dynamics
- neural networks: layers of nodes, connected by weight matrices
- weight matrices are updated using e.g. stochastic gradient descent (SGD)
- \circ stochastic matrix dynamics \rightarrow Dyson Brownian motion \rightarrow RMT features
- aim: further understanding of learning by characterising weight matrix dynamics
- identify universal behaviour and limitations of SGD during and after training

Outline

- o some general comments on stochastic weight matrix updates
- o connection to Dyson Brownian motion and stochastic Coulomb gas
- universal properties of stationary distribution
- application in Restricted Boltzmann Machine (RBM) and Transformer (nano-GPT)
- summary and outlook



Stochastic weight matrix dynamics

- $\,\circ\,\,$ consider some $M \times N\,\,$ weight matrix W
- update (e.g. stochastic gradient descent): $W \to W' = W + \delta W$ with $\delta W = -\alpha \frac{\delta \mathcal{L}}{\delta W}$
- $\,\circ\,\,$ obtained from loss function ${\cal L}[W]$, learning rate lpha
- δW is estimated using a batch \mathcal{B} with batch size $|\mathcal{B}|$: $\delta W_{\mathcal{B}} = \frac{1}{|\mathcal{B}|} \sum_{b \in \mathcal{B}} \delta W_b$
- fluctuations controlled by finite batch size (CLT): $\frac{1}{|\mathcal{B}|} \operatorname{Var}(\delta W)$

Stochastic weight matrix dynamics

 $\,\circ\,\,$ stochastic update $\,\,W \rightarrow W' = W + \delta W\,\,$ becomes

$$\delta W = \delta W_{\mathcal{B}} + \frac{1}{\sqrt{|\mathcal{B}|}} \sqrt{\operatorname{Var}(\delta W)} \eta$$

• or in terms of the gradient of the loss function:

$$W' = W - \alpha \left(\frac{\delta \mathcal{L}}{\delta W}\right)_{\mathcal{B}} + \frac{\alpha}{\sqrt{|\mathcal{B}|}} \sqrt{\operatorname{Var}\left(\frac{\delta \mathcal{L}}{\delta W}\right)} \eta \qquad \eta_{ij} \sim \mathcal{N}(0,$$

From rectangular to symmetric matrices

• W is $M \times N$ matrix: singular value decomposition: $W = U \Xi V^T$ $U U^T = \mathbb{1}$ $V V^T = \mathbb{1}$

○ singular values: ξ_i (i = 1...N) [take $N \le M$ without loss of generality]

o introduce symmetric semi-positive combination: $X = W^T W = V D V^T$

 \circ and focus on the singular/eigenvalues (invariant under left/right rotations on W):

$$D = \Xi^T \Xi = \operatorname{diag}\left(\xi_1^2, \dots, \xi_N^2\right) = \operatorname{diag}\left(x_1, \dots, x_N\right)$$

• stochastic dynamics: $X \to X' = X + \delta X_{\mathcal{B}} + \frac{1}{\sqrt{|\mathcal{B}|}} \sqrt{\operatorname{Var}(\delta X)} \eta$

Initialisation: Marchenko-Pastur distribution

o if initial weight matrix $W_{ij} \sim \mathcal{N}(0, \sigma^2)$ then X follows Marchenko-Pastur distribution

$$P_{\rm MP}(x) = \frac{1}{2\pi\sigma^2 Mrx} \sqrt{(x_+ - x)(x - x_-)} \qquad x_- < x < x_+ \qquad r = N/M \le 1 \quad x_\pm = M\sigma^2 \left(1 \pm \sqrt{r}\right)^2$$

how to choose σ^2 : distribution should depend on r only, safe to take large N, M limit

✓ spectrum is bounded for all r (relevant for RBMs below) : $\sigma^2 = 1/M$: $N \le M$

$$P_{\rm MP}(x) = \frac{1}{2\pi r x} \sqrt{(x_+ - x)(x - x_-)} \qquad 0 \le x_- \le x \le x_+ \le 4 \qquad x_\pm = \left(1 \pm \sqrt{r}\right)^2$$

Stochastic matrix dynamics: Dyson Brownian motion and the stochastic Coulomb gas

- $\,\circ\,\,$ framework to consider stochastic matrix dynamics for symmetric matrix X
- Dyson Brownian motion (in continuous time for now, see below):

$$\frac{dX_{ij}}{dt} = K_{ij}(X) + \sqrt{A_{ij}}\eta_{ij}$$

eigenvalues then evolve according to

$$egin{aligned} & rac{dx_i}{dt} = K_i(x_i) + \sum_{j
eq i} rac{g_i^2}{x_i - x_j} + \sqrt{2}g_i\eta_i \ & \equiv K_i^{(ext{eff})}(x_i) + \sqrt{2}g_i\eta_i \ & ext{where } \sqrt{A_{ii}} = \sqrt{2}g_i \end{aligned}$$

Dyson Brownian motion, stochastic Coulomb gas

• eigenvalues dynamics:
$$rac{dx_i}{dt} = K_i(x_i) + \sum_{j
eq i} rac{g_i^2}{x_i - x_j} + \sqrt{2}g_i\eta_i$$

- can be derived using 2nd order perturbation theory
- Coulomb term: eigenvalue repulsion [Wigner, Dyson 1959-1962, for nuclear spectra]
- Fokker-Planck equation (FPE) for distribution of eigenvalues:

$$\partial_t P(\{x_i\}, t) = \sum_{i=1}^N \partial_{x_i} \left[\left(g_i^2 \partial_{x_i} - K_i^{(\text{eff})}(x_i) \right) \right] P(\{x_i\}, t)$$

Dyson Brownian motion, stochastic Coulomb gas

• FPE:
$$\partial_t P(\{x_i\}, t) = \sum_{i=1}^N \partial_{x_i} \left[\left(g_i^2 \partial_{x_i} - K_i^{(\text{eff})}(x_i) \right) \right] P(\{x_i\}, t)$$

• stationary distribution: $P_s(\{x_i\}) = \frac{1}{Z} \prod_{i < j} |x_i - x_j| e^{-\sum_i V_i(x_i)/g_i^2}$
• with partition function: $Z = \int dx_1 \dots dx_N P_s(\{x_i\})$

o and provided drift can be derived from a potential $K_i(x_i) = -\frac{dV_i(x_i)}{dx_i}$

known as Coulomb gas, describes universal features of random matrices

Back to weight matrix dynamics

• stochastic dynamics
$$X \to X' = X + \delta X_{\mathcal{B}} + \frac{1}{\sqrt{|\mathcal{B}|}} \sqrt{\operatorname{Var}(\delta X)} \eta$$

• what can be carried over from Dyson's matrix dynamics? implications? universality?

$$\circ$$
 eigenvalue equation: $x_i \to x_i' = x_i + \delta x_i + \sum_{j \neq i} rac{g_i^2}{x_i - x_j} + \sqrt{2}g_i\eta_i$

• make explicit learning rate and batch size dependence

$$\delta x_i = \alpha K_i$$
 $g_i = \frac{\alpha}{\sqrt{|\mathcal{B}|}} \tilde{g}_i$ $\tilde{g}_i \sim \operatorname{Var}(\delta \mathcal{L}/\delta W)|_{ii}$

Back to weight matrix dynamics

$$\circ$$
 eigenvalue dynamics: $x_i o x_i' = x_i + \delta x_i + \sum_{j
eq i} rac{g_i^2}{x_i - x_j} + \sqrt{2} g_i \eta_i$

insert learning rate and batch size dependence:

$$x_i \to x'_i = x_i + \alpha K_i + \frac{\alpha^2}{|\mathcal{B}|} \sum_{j \neq i} \frac{\tilde{g}_i^2}{x_i - x_j} + \frac{\alpha}{\sqrt{|\mathcal{B}|}} \sqrt{2} \tilde{g}_i \eta_i$$

no usual scaling of drift and noise with learning rate (Ito calculus: ϵ , $\sqrt{\epsilon}$):
 no obvious continuous time limit (SDE), only in some weak sense

Q Li, C Tai and W E [1511.06251] S Yaida [1810.00004]

known issue: from SGD to SDE but is in fact blessing (see below)

$$x_i \to x_i' = x_i + \alpha K_i + \frac{\alpha^2}{|\mathcal{B}|} \sum_{j \neq i} \frac{\tilde{g}_i^2}{x_i - x_j} + \frac{\alpha}{\sqrt{|\mathcal{B}|}} \sqrt{2} \tilde{g}_i \eta_i$$

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Stationary distribution

- o distribution for fixed $lpha, |\mathcal{B}|$: $P_s(\{x_i\}) = rac{1}{Z} \prod_{i < j} |x_i x_j| e^{-\sum_i V_i(x_i)/g_i^2}$
- make explicit dependence on learning rate and batch size

$$g_i = rac{lpha}{\sqrt{|\mathcal{B}|}} ilde{g}_i \qquad \qquad V_i(x_i) = lpha ilde{V}_i(x_i) \qquad \qquad rac{V_i(x_i)}{g_i^2} = rac{1}{lpha/|\mathcal{B}|} rac{ ilde{V}_i(x_i)}{ ilde{g}_i^2}$$

• if drift vanishes at $x_i = x_i^s$, expand potential $\tilde{V}_i(x_i) = \tilde{V}_i(x_i^s) + \frac{1}{2}\Omega_i(x_i - x_i^s)^2 + \dots$

• exponential is Gaussian with variance $\sigma_i^2 = (\alpha/|\mathcal{B}|) (\tilde{g}_i^2/\Omega_i)$ universal scaling with model-dependent learning rate and batch size factor

Linear scaling relation

 \circ dependence on $\alpha/|\mathcal{B}|$ in training has been observed before, empirically

 ✓ P. Goyal, P. Dollár, R.B. Girshick, P. Noordhuis, L. Wesolowski, A. Kyrola et al., *Accurate, Large Minibatch SGD: Training ImageNet in 1 Hour* [1706.02677]
 ✓ S.L. Smith and Q.V. Le,

A Bayesian Perspective on Generalization and Stochastic Gradient Descent [1710.06451]

✓ S.L. Smith, P. Kindermans and Q.V. Le,

Don't Decay the Learning Rate, Increase the Batch Size [1711.00489]

o finds a natural place in the framework of Dyson Brownian motion and Coulomb gas

Applications and implications

o so far, the derivation is general: prediction of eigenvalue distribution after learning

- o apply to actual ML models to observe universal features and support derivation
- teacher-student model
- Gaussian Restricted Boltzmann Machine
- Transformer

builds on previous analysis of RBM: GA, B Lucini, C Park, Phys. Rev. D 109 (2024) 034521 [2309.15002 [hep-lat]] current analysis: PRE 111 (2025) 1, 015303 [2407.16427 [cond-mat.dis-nn]]

GA, O Hajizadeh, B Lucini, C Park 2411.13512 [cond-mat.dis-nn]

Restricted Boltzmann Machine: generative network



- energy-based method
- probability distribution
- binary or continuous d.o.f.

$$p(\phi,h) = rac{1}{Z} e^{-S(\phi,h)}$$

$$Z = \int D\phi Dh \, e^{-S(\phi,h)} \, \, {}_{\rm 18}$$



Scalar field RBM

o distribution:
$$p(\phi,h) = \frac{1}{Z}e^{-S(\phi,h)}$$
 $S(\phi,h) = \frac{1}{2}\mu^2\phi^T\phi + \frac{1}{2\sigma_h^2}(h-\eta)^T(h-\eta) - \phi^TWh$

- $\circ M imes N = N_v imes N_h$ weight matrix W
- induced distribution on visible layer $p(\phi) = \int Dh \, p(\phi, h) = \frac{1}{Z} \exp\left(-\frac{1}{2}\phi^T K \phi + J^T \phi\right)$
- $\circ \quad \text{kernel} \quad K = \mu^2 \mathbbm{1} \sigma_h^2 W W^T = \mu^2 \mathbbm{1} \sigma_h^2 U \Xi \Xi^T U^T = U \left[\mu^2 \mathbbm{1} \sigma_h^2 \Xi \Xi^T \right] U^T \equiv U D_K U^T$

• eigenvalues
$$D_K = \operatorname{diag}\left(\underbrace{\mu^2 - \sigma_h^2 \xi_1^2, \mu^2 - \sigma_h^2 \xi_2^2, \dots, \mu^2 - \sigma_h^2 \xi_N^2}_{N}, \underbrace{\mu^2, \dots, \mu^2}_{M-N}\right)$$

Scalar field RBM as a lattice field theory

• treat RBM as a lattice field theory with bi-linear quadratic action

$$S(\phi, h) = \sum_{i} \frac{1}{2} \mu_{i}^{2} \phi_{i}^{2} + \sum_{a} \frac{1}{2\sigma^{2}} (h_{a} - \eta_{a})^{2} - \sum_{i,a} \phi_{i} w_{ia} h_{a}$$

induced distribution on visible layer

$$p(\phi) = \int Dh \, p(\phi, h) = \frac{1}{Z} \exp\left(-\frac{1}{2} \sum_{i,j} \phi_i K_{ij} \phi_j + \sum_i J_i \phi_i\right)$$

all information is stored in quadratic operator, with spectrum

$$D_{K} = \operatorname{diag}\left(\underbrace{\mu^{2} - \sigma_{h}^{2}\xi_{1}^{2}, \mu^{2} - \sigma_{h}^{2}\xi_{2}^{2}, \dots, \mu^{2} - \sigma_{h}^{2}\xi_{N}^{2}}_{N}, \underbrace{\mu^{2}, \dots, \mu^{2}}_{M-N}\right)$$

Scalar field RBM as an ultraviolet regulator

○ spectrum

$$D_{K} = \operatorname{diag}\left(\underbrace{\mu^{2} - \sigma_{h}^{2}\xi_{1}^{2}, \mu^{2} - \sigma_{h}^{2}\xi_{2}^{2}, \dots, \mu^{2} - \sigma_{h}^{2}\xi_{N_{h}}^{2}}_{N_{h}}, \underbrace{\mu^{2}, \dots, \mu^{2}}_{N_{v} - N_{h}}\right)$$

- \circ what if $N_h < N_v$? not all eigenvalues can be reproduced
- $_{\odot}$ role of hyperparameter μ^2 ? if chosen too low, not all eigenvalues can be reproduced

$$lacksim$$
 both N_h and μ^2 act as ultraviolet regulators

GA, B Lucini, C Park, Phys. Rev. D 109 (2024) 034521 [2309.15002 [hep-lat]]

RBM as ultraviolet regulator



- apply to MNIST data set (28 x 28 images)
- compute spectrum of two-point
 correlator $K_{ij}^{-1} = \langle \phi_i \phi_j
 angle_{
 m data}$
- \circ inverse spectrum $1/\kappa$
- infrared safe
- ultraviolet divergent

C	infrared	 6.572
0		 4.806
		 4.178
		 3.650
, ک		 3.297
-		 2.920 -
-		 2.216
2		 1.953
		 1.871
		 1.596
0	ultraviolet	1.428

784 eigenvalues



MNIST with fixed RBM mass

- $\circ \ N_v = N_h = 784$
- $\circ~$ fixed RBM mass $\mu^2=100$
- spectrum regulated
- infrared modes learned approximately correctly (see below)



MNIST with $N_h \leq N_v$

what is the effect of including incomplete spectrum?

5	0	Ч	1	9	2	١	3
1	4	3	ک	3	6	1	7
9	8	6	9	T	0	9	1
ユ	г	4	3	2	7	N	8

5	0	Ч	1	9	2	١	3
1	4	3	ک	3	6	1	7
Υ	8	6	9	T	0	9	1
ユ	г	Ч	3	2	7	Ы	8



removal of

ultraviolet modes

affects

generative power

(a) $N_h = 784$

5	0	H	1	9	3	1	З
1	4	3	Ś	3	6	Ŧ	7
Э	8	6	9	ч	0	9	1
<u>t</u>	З	4	ß	2	7	3	8

(d) $N_h = 36$

(b) $N_h = 225$

(c) $N_h = 64$

53	0	q	1	9	3	4	З
3	9	3	\mathcal{G}	3	6	4	7
3	8	6	9	5	0	9	1
(2)	3	4	G	3	4	3	8

(e) $N_h = 16$



(f) $N_h = 4$

Back to Dyson Brownian motion

- weight matrix is updated using persistent contrastive divergence (PCD)
- o maximise likelihood/minimise KL divergence

$$\frac{\delta \mathcal{L}}{\delta W_{ia}} = \sigma_h^2 \big(\langle \phi_i \phi_j \rangle_{\text{target}} - \langle \phi_i \phi_j \rangle_{\text{model}} \big) W_{ja}$$

- o denote eigenvalues of $X = W^T W$ as x_i
- PCD is stochastic:

$$x_i \to x'_i = x_i + \alpha K_i + \frac{\alpha^2}{|\mathcal{B}|} \sum_{j \neq i} \frac{\tilde{g}_i^2}{x_i - x_j} + \frac{\alpha}{\sqrt{|\mathcal{B}|}} \sqrt{2} \tilde{g}_i \eta_i$$

 N_h nodes

 $N_{\rm v}$ nodes

W

Back to Dyson Brownian motion

o maximise likelihood/minimise KL divergence

$$\frac{\delta \mathcal{L}}{\delta W_{ia}} = \sigma_h^2 \big(\langle \phi_i \phi_j \rangle_{\text{target}} - \langle \phi_i \phi_j \rangle_{\text{model}} \big) W_{ja}$$

- \circ denote target distribution has eigenvalues with κ_i
- drift in instantaneous eigen-basis: $K_i(x_i) = \left(\frac{1}{\kappa_i} \frac{1}{\mu^2 x_i}\right) x_i$
- \circ fixed point of drift: $x_i^s = \mu^2 \kappa_i$, spectrum learnt correctly
- o where can we observe the effects of RMT?

Scalar field RBM

- o implement for simple target distribution: scalar field in LFT in 1 dimension
- spectrum is free dispersion relation: $\kappa_k = m^2 + p_{\text{lat},k}^2 = m^2 + 2 2\cos\left(\frac{2\pi k}{N_v}\right)$
- each mode is doubly degenerate, except lowest and highest one
- example for 10 modes
 degenerate modes split for clarity
 4

RBM evolution

weight matrix updates using persistent contrastive divergence with mini-batches



initial Marchenko-Pastur distribution

towards target spectrum

RBM evolution and RMT universality

- weight matrix updates using persistent contrastive divergence with mini-batches
- no sharp lines, distributions around target spectrum
- test predictions from RMT:
 - induced Coulomb term and eigenvalue repulsion (universal)
 - potential from drift (non-universal)



Universal RMT predictions

consider two degenerate modes only: Coulomb gas description

$$Z = \frac{1}{N_0} \int dx_1 dx_2 |x_1 - x_2| e^{-V(x_1, x_2)} \qquad V(x_1, x_2) = \frac{1}{2\sigma^2} \left[(x_1 - \kappa)^2 + (x_2 - \kappa)^2 \right]$$

 \circ eigenvalues x_1, x_2 cannot both be equal to κ due to Coulomb repulsion

• two ways to detect this: Wigner surmise and Wigner semi-circle

• Wigner surmise: distribution for level spacing $S = x_1 - x_2$

$$P(S) = \frac{S}{2\sigma^2} e^{-S^2/(4\sigma^2)}.$$

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• mean level spacing
$$\langle S \rangle = \int_0^\infty dS \, SP(S) = \sqrt{\pi} \sigma$$
. $s = S/\langle S \rangle$

Wigner surmise

$$\odot$$
 distribution $P(S) = rac{S}{2\sigma^2} e^{-S^2/(4\sigma^2)}$ for level spacing $S = x_1 - x_2^{0.0}$

• mean level spacing
$$\langle S \rangle = \int_0^\infty dS \, SP(S) = \sqrt{\pi} \sigma$$
.

$$\circ$$
 Wigner surmise for $s=S/\langle S
angle$: $P(s)=rac{\pi}{2}se^{-\pi s^2/4}$ universal curve

- many RBM training runs, stochasticity due to mini-batches, collect histograms of x_i
- vary learning rate and batch size [no ordering of eigenvalues by hand, induces bias!]

2

S

3

0.8

0.6

୍ଡ ଜୁ 0.4

0.2



Wigner surmise: 4 degenerate pairs

$$P(S) = \frac{S}{2\sigma^2} e^{-S^2/(4\sigma^2)}, \qquad \langle S \rangle = \sqrt{\pi}\sigma \qquad P(s) = \frac{\pi}{2} s e^{-\pi s^2/4}$$



Wigner surmise: vary learning rate and batch size

• prediction:

 $\sigma_i^2 = (lpha / |\mathcal{B}|) \left(ilde{g}_i^2 / \Omega_i
ight)$

- \circ linear dependence on $(lpha/|\mathcal{B}|)$
- mean level spacing

$$egin{aligned} \langle S_i
angle &= \pi \sqrt{(lpha / |\mathcal{B}|)(ilde{g}_i^2 / \Omega_i)} \ &= a_{ ext{fit}} \sqrt{(lpha / |\mathcal{B}|)(\kappa_i^2 \Omega_i)} \end{aligned}$$

fit function includes

non-universal parameters as well





• for two modes: $\rho(x) = \frac{e^{-x^2/(2\sigma^2)}}{4\sqrt{\pi}\sigma} \left[2e^{-x^2/(2\sigma^2)} + \sqrt{2\pi}\frac{x}{\sigma} \operatorname{Erf}\left(\frac{x}{\sqrt{2}\sigma}\right) \right]$

- broadened and flattened Gaussian
- fit σ parameter and position for each doubly degenerate mode

Wigner semi-circle

o fit to semi-circle for two different κ_i values with fixed learning rate and batch size

 $\rho(x) = \frac{e^{-x^2/(2\sigma^2)}}{4\sqrt{\pi}\sigma} \left[2e^{-x^2/(2\sigma^2)} + \sqrt{2\pi}\frac{x}{\sigma} \operatorname{Erf}\left(\frac{x}{\sqrt{2}\sigma}\right) \right]$

• Binder cumulant $U_4 = -4/27 \approx -0.148$ for semi-circle (vanishes for Gaussian)



Wigner semi-circle and surmise

semi-circle

dependence on learning rate/batch size



consistency between surmise and semi-circle fits



Wigner surmise and semi-circle

✓ parameter σ scales as: $\sigma_i^2 = (\alpha/|\mathcal{B}|) \quad (\tilde{g}_i^2/\Omega_i)$ universal scaling model-dependent

✓ stochasticity leads to universal features in trained models

derived that learning rate and finite batch size appear as ratio

✓ previously observed as empirical linear scaling rule

Eigenvalue repulsion

- Coulomb interaction between all eigenvalues
- learned eigenvalue/target
- repulsion for nonzero learning rate/batch size
- no "perfect learning" unless stochasticity vanishes
- overfitting, generalisation, ...



Non-universal dynamics

• consider this for one mode only (drop the index)

$$V(x) = -\int^{x} dx' K(x') = -\frac{x^{2}}{2\kappa} - x - \mu^{2} \log(\mu^{2} - x)$$

• stationary distribution

$$P_s(x) = \frac{1}{Z} e^{-V(x)/g^2} = \frac{1}{Z} \exp\left[\frac{1}{g^2} \left(\frac{x^2}{2\kappa} + x + \mu^2 \log\left(\mu^2 - x\right)\right)\right]$$

Time-dependent dynamics

- assume continuous time limit exists
- analyse FPE for one mode: $\partial_{\tau} P(x,\tau) = \partial_x \left(g^2 \partial_x K(x) \right) P(x,\tau)$
- o solve using standard stochastic quantisation/FP methods: $P(x, \tau) = \sqrt{P_s(x)}\psi(x, \tau)$

• evolution:
$$\partial_{\tau}\psi(x,\tau) = \left(g^2\partial_x^2 - \frac{1}{4g^2}\left[\partial_x V(x)\right]^2 + \frac{1}{2}\left[\partial_x^2 V(x)\right]\right)\psi(x,\tau) \equiv -2H_{\rm FP}\psi(x,\tau)$$

• Fokker-Planck Hamiltonian: $H_{\rm FP} = \frac{1}{2}L^{\dagger}L$
 $L^{\dagger} = -g\partial_x + \frac{1}{2g}\partial_x V(x)$
 $L = +g\partial_x + \frac{1}{2g}\partial_x V(x)$ 40

Quantum-mechanical bound state problem

• $H_{\rm FP} = \frac{1}{2}L^{\dagger}L$ eigenvalue problem: $H_{\rm FP}\psi_n(x) = E_n\psi_n(x)$



$$U(x) = \frac{1}{g^2} \left[U_0(x) + g^2 U_1(x) \right]$$
$$U_0(x) = \frac{1}{8} \left[\partial_x V(x) \right]^2$$
$$U_1(x) = -\frac{1}{4} \partial_x^2 V(x)$$

double well potential on interval $0 \le x \le \mu^2$

Quantum-mechanical bound state problem

• $H_{\rm FP}\psi_n(x) = E_n\psi_n(x)$ ground state exactly known: $\psi_0(x) = \sqrt{P_s(x)}$

 \circ width of solution depends on strength of the noise g^2 : better description of target κ



Full time-dependent dynamics: learning

- combine Coulomb repulsion and drift
- from Marchenko-Pastur distribution
 to stochastic target distribution
- 10 modes, 4 doubly degenerate ones
- dynamics of $P(\{x_i\}, t)$ described by FPE



effective description of learning dynamics in terms of eigenvalues

Second application: Transformers

• Gaussian RBM has one weight matrix, target spectrum is known, essentially solvable

in more advanced architectures:

• many weight matrices, target spectra not known, do they even exist?

• what is the loss function landscape? localised minima, flat directions, ... ?

• empirical study following dynamics of eigenvalues of $X = W^T W$

GA, O Hajizadeh, B Lucini, C Park, NeurIPS 2024 workshop *ML and the Physical Sciences*, <u>2411.13512</u> [cond-mat.dis-nn]

Transformer: nano-GPT

- four attention blocks with each four attention heads
- \circ each attention head: one key (K), one query (Q) and one value (V) matrix
- matrix sizes: $M \times N = 64 \times 16$
- \circ about 2.1 \times 10⁵ parameters
- use AdamW optimiser
 - (highly adaptive stepsize during training)
- trained on opus of Shakespeare



• initialisation: eigenvalues of $X = W^T W$ follow Marchenko-Pastur distribution



- \circ appearance of tail in distribution (shown K matrix of layer 1)
- part of spectrum described by Marchenko-Pastur distribution is reduced, A < 1
- \circ use area A and width σ^2 as fit parameters during evolution



 \circ evolution of area A and width σ^2 during evolution (shown K matrix for all four layers)





15-25% of spectral weight moves to the tail

MP distribution broadens due to Brownian motion

Transformer: Wigner surmise

- short-distance fluctuations: level spacing described by Wigner surmise
- o remains approximately described by RMT for real, symmetric matrices



iteration 0

iteration 1000

iteration 5000⁴⁹

requires further understanding:

- what is the "final" target spectrum? does it even exist?
- tail drops as a power, what does this imply? can the power be understood?

significant part of the spectrum remains MP: random matrix elements

o how relevant is this part of the spectrum? remove? sparse weight matrices?

see also CH Martin, MW Mahoney, Traditional and Heavy-Tailed Self Regularization in Neural Network Models, 1901.08276



- o stochastic weight matrix dynamics has universal features described by RMT
- o manifests in eigenvalue repulsion, quantified by Wigner surmise and semi-circle
- o fundamental limitation of learning for finite learning rate and batch size
- stochasticity controlled by learning rate/batch size: reduce ratio to improve agreement with target distribution, but stochasticity allows for generalisation

Outlook

- Dyson Brownian motion is present at "microscopic" level
- o how does it manifest itself for more advanced architectures?
- is there universality beyond level repulsion (power law tails)?
- what are the practical implications? description of learning, algorithmic advances?